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Date: May 24, 2022.

Abstract.

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1. INTRODUCTION

1.1. Dirac Delta Function [1]. The dirac delta function δ is a generalized function over the real numbers where it is zero everywhere except at 0. We can construct it using a limit of the *delta*-sequence given below,

$$\delta_n(t) = \begin{cases} \frac{n}{2} & -\frac{1}{n} < t < \frac{1}{n} \\ 0 & \text{otherwise} \end{cases}$$

From the above sequence we can see that $\delta_n(t)$ converges to a spike (as $n \to \infty$) at t = 0 with infinite height and zero width with the property

$$\int_{-\infty}^{\infty} \delta_n(t) dt = 0, \forall n \ge 1$$

Now assume a continuous function f(t). We can write

$$\min_{t \in \left[\frac{-1}{n}, \frac{1}{n}\right]} f(t) \le \int_{-\infty}^{\infty} f(t) \delta_n(t) dt \le \max_{t \in \left[\frac{-1}{n}, \frac{1}{n}\right]} f(t)$$

Since,

$$\lim_{n \to \infty} \min_{t \in \left[\frac{-1}{n}, \frac{1}{n}\right]} f(t) = \lim_{n \to \infty} \max_{t \in \left[\frac{-1}{n}, \frac{1}{n}\right]} f(t) = f(0)$$

then by the squeeze theorem, we can write

$$\lim_{n \to \infty} \int_{-\infty}^{\infty} f(t) \delta_n(t) dt = f(0).$$

Definition 1. The dirac delta function can be defined by the following equality:

$$\int_{-a}^{a} f(t)\delta(t)dt = f(0)$$

for any a > 0 and any continuous function f. This gives us the following relation,

$$\int_{s-a}^{s+a} f(t)\delta(t-s)d\tau = f(s)$$

(This is also called the Sifting property.)

1.2. Spectral Theory.

Definition 2. (pre-Hilbert) Let X be a complex vector space. A Hermitian inner product on X is a function $\langle \cdot, \cdot \rangle : X \times X \longrightarrow \mathbb{C}$, which is :

(1) (positive non-degenerate) $(\forall x \in X) \langle x, x \rangle \ge 0, \langle x, x \rangle = 0$ iff x = 0.

 $(2) \ (sesquilinear) \ (\forall \alpha, \beta \in \mathbb{C}) (\forall x, y, z \in X), \langle \alpha x + \beta y, z \rangle = \alpha \langle x, z \rangle + \beta \langle y, z \rangle$

(3) (conjugate-symmetric) $(\forall x, y \in X) \langle y, x \rangle = \overline{\langle x, y \rangle}$

A space is complete if all cauchy sequences converge to a limit . A *Hilbert Space* is a complete inner product space. In fact every Hilbert space is a *Banach Space* but the reverse is not true.

Definition 3. Banach space: Let V be a vector space. A norm is a mapping $\|\cdot\| : V \to [0, \infty)$ that satisfies:

- (1) $||x + y|| \le ||x|| + ||y||.$
- (2) ||ax|| = a||x|| for all $a \in \mathbb{R}$.
- (3) ||x|| = 0 implies that x = 0.

A complete normed space is called a Banach space.

We can now define an inner product in function space. An inner product in the vector space of continuous functions denoted by $\mathcal{C}^0([a,b],\mathbb{C})$ is defined as follows, given two arbitrary functions $u(x), v(x), x \in [a,b]$, introduce the inner product,

$$\langle u, v \rangle = \int_{a}^{b} uv^{*} dx, \quad *: \text{ complex conjugation.}$$

Theorem 1. (Fredholm Alternative theorem)

For a linear system Ax = b, A^* which is the adjoint satisfying the property $\langle Ax, y \rangle = \langle x, A^*y \rangle$ for all $x, y, \in \mathbb{C}^n$ exactly one is true:

i) Solution of Ax = b, if it exists, is unique if and only if x = 0 is the only solution of Ax = 0.

ii) The equation Ax = b has a solution if and only if $\langle b, v \rangle = 0$ for all v satisfying $A^*v = 0$.

Proof: i) Assume that Ax = 0 for $x \neq 0$ and $Ax_0 = b$. Then $A(x_0 + \alpha x) = b$ for all α . Therefore, the solution is not unique. Conversely, if there are two different solutions, x_1 and x_2 , satisfying $Ax_1 = b$ and $Ax_2 = b$, then one has a nonzero solution $x = x_1 - x_2$ such that $Ax = A(x_1 - x_2) = 0$. ii) Let $A^*v = 0$ and $Ax_0 = b$. Then we have

$$\langle b, v \rangle = \langle Ax_0, v \rangle = \langle x_0, A^*v \rangle = 0$$

For the second part we assume that $\langle b, v \rangle = 0$ for all v such that $A^*v = 0$. Write b as the sum of a part that is in the range of A and a part that in the space orthogonal to the range of $A, b = b_r + b_o$. Then, $0 = \langle b_o, Ax \rangle = \langle A^*b, x \rangle$ for all x. Since $\langle b, v \rangle = 0$ for all v in the nullspace of A^* , then $\langle b, b_o \rangle = 0$ Therefore, $\langle b, v \rangle = 0$ implies that

$$0 = \langle b, b_o \rangle = \langle b_r + b_o, b_o \rangle = \langle b_o, b_o \rangle$$

This means that $b_o = 0$, giving $b = b_r$ is in the range of A. So, Ax = b has a solution.

The same can be said about linear operators, let L be a bounded linear operator on a Hilbert space with adjoint L^{\dagger} . Then exactly one of the following is true:

i) The inhomogeneous problem

Lu = f

 $L^{\dagger}u = 0$

has a unique solution u.ii) The homogeneous adjoint problem

has a non-trivial solution.

Fredholm Alternative theorem is established by taking the inner product of (1.1) with the adjoint null space function v(x) and we obtain,

1.3. Green's Functions. The focus of this paper is to construct the appropriate Green's function for the following BVP (A):

$$\begin{cases} x'' + \lambda^2 x = \delta(t-s) \\ x(0) = x(\omega) \\ x'(0) = x'(\omega) \end{cases}$$

We first solve Ly = f, a differential equation with homogenous boundary conditions. The Sifting property mentioned earlier can be written as,

$$\langle f(t), \delta(t-s) \rangle = \int_0^l f(t)\delta(t-s) dt = f(s), \quad t \in [0, l]$$

Now we can consider the following problem,

$$Lu = f$$

 $L^{\dagger}G(t,s) = \delta(t-s)$

where $x \in [0, l]$ and L^{\dagger} is the adjoint operator with its boundary conditions. G(t, s) is the Greens function. Taking the inner product of Lu = f with respect to G(t, s) we get,

This gives us

$$u(t) = \int_0^l f(t)G(t,s)dt$$

and thus the inverse operator

$$L^{-1}[f] = \int_0^l f(t)G(t,s)ds$$

can be computed.

2. Solution of
$$x'' + \lambda^2 x = \delta(t-s)$$

From the above section we can define the Greens function G(t, s) of \mathcal{L} to be the unique solution to the problem $\mathcal{L}G = \delta(t-s)$, where \mathcal{L} is the general linear second order differential operator. We now construct the green function for A: From the definition of the dirac delta function, at all points $t \neq s$, $x'' + \lambda^2 x = 0$. Solving this we get homogenous solution is $k \cos(\lambda t) + c \sin(\lambda t)$. Therefore our Greens function is,

(1)
$$G(t,s) = \begin{cases} c_1 \cos(\lambda t) + c_2 \sin(\lambda t) & 0 \le t < s \le \omega \\ c_3 \cos(\lambda t) + c_4 \sin(\lambda t) & 0 \le s < t \le \omega \end{cases}$$

To solve for G(t,s) we have to solve for c_1, c_2, c_3, c_4 , we can do this by solving the following system of equations,

The periodic boundary conditions,

(2)
$$x(0) = x(\omega), \ k\cos(0) + c\sin(0) = k\cos(\lambda\omega) + c\sin(\lambda\omega)$$

(3)
$$x'(0) = x'(\omega), \ -k\lambda\sin(0) + c\lambda\cos(0) = -k\lambda\sin(\lambda\omega) + c\lambda\cos(\lambda\omega)$$

At t = s (as G(t, s) is continuous),

$$\lim_{t\to s^+} G(t,s) - \lim_{t\to s^-} G(t,s) = 0$$

(4)
$$c_3\cos(\lambda s) + c_4\sin(\lambda s) - c_1\cos(\lambda s) - c_2\sin(\lambda s) = 0$$

and the "derivative" jump of G(t, s),

$$\lim_{t \to s^+} G'(t,s) - \lim_{t \to s^-} G'(t,s) = 1$$

(5)
$$-c_3\lambda\sin(\lambda s) + c_4\lambda\cos(\lambda s) - (-c_1\lambda\sin(\lambda s) + c_2\lambda\cos(\lambda s)) = 1$$

Remark 1. Why $\lim_{t\to s^+} G'(t,s) - \lim_{t\to s^-} G'(t,s) = 1$?

Consider $\mathcal{L}G = \delta(t-s)$, this is zero when t < s and t > s and δ is infinity when t = s. Which tells us the first derivative must be discontinuous and when we take the second derivative it diverges. Now integrate the given ode from $s - \epsilon$ to $s + \epsilon$ and let $\epsilon \to 0$. We get,

$$\int_{s-\epsilon}^{s+\epsilon} \frac{\partial G}{\partial t^2} dt + \int_{s-\epsilon}^{s+\epsilon} \lambda^2 G = \int_{s-\epsilon}^{s+\epsilon} \delta(t-s) dt = 1$$

The second term on the l.h.s vanishes as $\epsilon \to 0$ as the integrands are finite and so we get,

$$\left. \frac{\partial G}{\partial t} \right|_{t=s^+} - \left. \frac{\partial G}{\partial t} \right|_{t=s^-} = 1$$

Solving (2), (3), (4) and (5) we get our Greens Function,

(6)
$$G(t,s) = \begin{cases} \frac{\sin(\lambda(t-s+\omega)) + \sin(\lambda(s-t))}{2\lambda(1-\cos(\lambda\omega))} & 0 \le t < s \le \omega\\ \frac{\sin(\lambda(s-t+\omega)) + \sin(\lambda(t-s))}{2\lambda(1-\cos(\lambda\omega))} & 0 \le s < t \le \omega \end{cases}$$

(7)
$$G(t,s) = \begin{cases} G_1(t,s) & 0 \le t < s \le \omega \\ G_2(t,s) & 0 \le s < t \le \omega \end{cases}$$

We can verify this as,

$$\lim_{t \to s^{+}} G(t, s) - \lim_{t \to s^{-}} G(t, s) = 0$$
$$\lim_{t \to s^{+}} G'(t, s) - \lim_{t \to s^{-}} G'(t, s) = 1$$
$$G(t, s) = G(s, t)$$

3. Solution of x'' = f(t, x(t), x'(t))

Consider the following B.V.P (B),

$$\left\{ \begin{array}{l} x^{\prime\prime}=f(t,x(t),x^{\prime}(t))\\ x(0)=x(\omega)\\ x^{\prime}(0)=x^{\prime}(\omega) \end{array} \right.$$

Notice the above ODE can be written as,

(8)
$$x'' + \lambda^2 x(t) = \lambda^2 x(t) + f(t, x(t), x'(t))$$

Note: The solution to the differential equation $x'' + {}^2 x = \delta(t - s)$ is the Greens function we obtained above. Lemma 1. I claim the solution to the above differential equation (20) is,

(9)
$$x(t) = \int_0^\omega G(t,s) \left\{ f(s,x(s),x'(s)) + \lambda^2 x(s) \right\} ds$$

Why so? Existence discussed later

Proof: Need to show x(t) satisfies the BVP (B),

$$\begin{aligned} x(t) &= \int_{0}^{\omega} G(t,s) \left\{ f\left(s, x(s), x'(s)\right) + \lambda^{2} x(s) \right\} ds \\ x'(t) &= \frac{d}{dt} \int_{0}^{\omega} G(t,s) h(s) ds, \quad h(s) = \left\{ f\left(s, x(s), x'(s)\right) + \lambda^{2} x(s) \right\} \\ &= \frac{d}{dt} \int_{0}^{t^{-}} G_{2}(t,s) h(s) ds + \frac{d}{dt} \int_{t^{+}}^{\omega} G_{1}(t,s) h(s) ds \\ &= G_{2}(t,t^{-}) h(t^{-}) + \int_{0}^{t^{-}} \frac{\partial G_{2}(t,s)}{\partial t} h(s) ds - G_{1}(t,t^{+}) h(t^{+}) + \frac{d}{dt} \int_{t^{+}}^{\omega} \frac{\partial G_{1}(t,s)}{\partial t} h(s) ds \\ &= \int_{0}^{t^{-}} \frac{\partial G_{2}(t,s)}{\partial t} h(s) ds + \int_{t^{+}}^{\omega} \frac{\partial G_{1}(t,s)}{\partial t} h(s) ds \end{aligned}$$

$$\begin{aligned} x''(t) &= \frac{\partial G_2(t,t^-)}{\partial t} h(t^-) + \int_0^{t^-} \frac{\partial^2 G_2(t,s)}{\partial t^2} h(s) ds - \frac{\partial G_1(t,t^+)}{\partial t} h(t^+) + \int_{t^+}^{\omega} \frac{\partial^2 G_2(t,s)}{\partial t^2} h(s) ds \\ &= h(t) + \int_0^{t^-} \frac{\partial^2 G_2(t,s)}{\partial t^2} h(s) ds + \int_{t^+}^{\omega} \frac{\partial^2 G_1(t,s)}{\partial t^2} h(s) ds \\ &= h(t) + \lambda^2 x(t) - \lambda^2 x(t) + \int_0^{t^-} \frac{\partial^2 G_2(t,s)}{\partial t^2} h(s) ds + \int_{t^+}^{\omega} \frac{\partial^2 G_1(t,s)}{\partial t^2} h(s) ds - \lambda^2 x(t) \\ &= h(t) + \lambda^2 x(t) + \int_0^{t^-} \frac{\partial^2 G_2(t,s)}{\partial t^2} h(s) ds + \int_{t^+}^{\omega} G_1(t,s) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \lambda^2 \left(\int_0^{t^-} G_2(t,s) h(s) ds + \int_{t^+}^{\omega} G_1(t,s) h(s) ds \right) \right) + \\ \int_0^{t^-} \frac{\partial^2 G_2(t,s)}{\partial t^2} h(s) ds + \int_{t^+}^{\omega} \frac{\partial^2 G_1(t,s)}{\partial t^2} h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_0^{t^-} \left(\frac{\partial^2 G_2(t,s)}{\partial t^2} + \lambda^2 G_2(t,s) \right) h(s) ds + \\ \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_0^{t^-} \left(\frac{\partial^2 G_2(t,s)}{\partial t^2} + \lambda^2 G_2(t,s) \right) h(s) ds + \\ \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_0^{t^-} \left(\frac{\partial^2 G_2(t,s)}{\partial t^2} + \lambda^2 G_2(t,s) \right) h(s) ds + \\ \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_0^{t^-} \left(\frac{\partial^2 G_2(t,s)}{\partial t^2} + \lambda^2 G_2(t,s) \right) h(s) ds + \\ \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_0^{t^-} \left(\frac{\partial^2 G_2(t,s)}{\partial t^2} + \lambda^2 G_2(t,s) \right) h(s) ds + \\ \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_0^{t^-} \left(\frac{\partial^2 G_2(t,s)}{\partial t^2} + \lambda^2 G_2(t,s) \right) h(s) ds + \\ \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_0^{t^-} \left(\frac{\partial^2 G_2(t,s)}{\partial t^2} + \lambda^2 G_2(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1(t,s) \right) h(s) ds - \lambda^2 x(t) \\ &= h(t) + \int_{t^+}^{\omega} \left(\frac{\partial^2 G_1(t,s)}{\partial t^2} + G_1($$

 $x'' = h(t) - \lambda^2 x(t) algorithmic$ Rearranging we get,

$$x^{\prime\prime}+\lambda^2 x(t)=\lambda^2 x(t)+f(t,x(t),x^\prime(t))\implies x^{\prime\prime}=f(t,x(t),x^\prime(t))$$

Now we have to show that the boundary conditions are satisfied, i.e

$$x(0) = x(\omega)$$
$$x'(0) = x'(\omega)$$

Consider $x(0) = x(\omega)$,

(10)
$$x(t) = \int_0^{t^-} G_2(t,s)h(s)ds + \int_{t^+}^{\omega} G_1(t,s)h(s)ds$$
$$x(0) = \int_0^{\omega} G_1(0,s)h(s)ds$$
$$x(\omega) = \int_0^{\omega} G_2(\omega,s)h(s)ds$$

Well $x(0) = x(\omega)$ as we know G(t, s) satisfies the BVP (A). Another way to verify is by just evaluating $G_1(0, s)$ and $G_2(\omega, s)$. Consider $x'(0) = x'(\omega)$,

(11)
$$x'(t) = \int_{0}^{t^{-}} \frac{\partial G_{2}(t,s)}{\partial t} h(s) ds + \int_{t^{+}}^{\omega} \frac{\partial G_{1}(t,s)}{\partial t} h(s) ds$$
$$x'(0) = \int_{0}^{\omega} \frac{\partial G_{1}(0,s)}{\partial t} h(s) ds$$
$$x'(\omega) = \int_{0}^{\omega} \frac{\partial G_{2}(\omega,s)}{\partial t} h(s) ds$$

By a similar argument we get $x'(0) = x'(\omega)$. Therefore we showed the integral satisfies the ODE.

4. Degree Theory

Let A be a mapping: $A : \overline{U} \subset \mathbb{R}^N \to \mathbb{R}^N$, where U is an open bounded set. We defined the degree of the mapping at p,

$$\deg(A, U, p) = \sum_{x_i \in A^{-1}(p)} \operatorname{sign} \left(\det J_A(x_i) \right)$$

as long as $A^{-1}(p)$ are regular points. Let $h: \mathbb{R}^N \to \mathbb{R}$ is a smooth function such that

$$\int_{\mathbb{R}^N} h(x) dx = 1,$$

and h(x) = 0 outside of a ball $B_{\varepsilon}(0)$ for some small $\varepsilon > 0$. Here $x = (x_1, \ldots, x_N)$ and $dx = dx_1 \cdots dx_N$.

$$\deg(A, U, 0) = \int_{\mathbb{R}^N} h(A(x)) \det J_A(x) dx$$

for $x = (x_1, ..., x_n)$.

Lemma 2. We show the above integral is independent of h.

Proof. By induction. For n = 1, Let $\eta(x)$ be another function with support in $(-\varepsilon, \varepsilon)$ and

$$\int_{\mathbb{R}} \eta(x) dx = 1$$

Therefore, $\omega = h(x) - \eta(x)$ has the property

$$\int_{\mathbb{R}} \omega(x) dx = 0$$

We show

$$\int_{\mathbb{R}} \{h(A(x)) - \eta(A(x))\} A'(x) dx = 0$$

For simplicity, let us denote $h(A(x)) - \eta(A(x)) A'(x) dx = f(x) dx$. We show, there is g with support in $(-\varepsilon, \varepsilon)$ such that

$$f(x)dx = d(g)(x)$$

It is enough to take g as

$$g(x) = \int_{-\infty}^{x} f(x) dx$$

$$\int_{\mathbb{R}} f(x)dx = \int_{\mathbb{R}} d(g)(x) = g(\infty) - g(-\infty) = 0$$

and therefore

$$\int_{\mathbb{R}} h(A(x))A'(x)dx = \int_{\mathbb{R}} \eta(A(x))A'(x)dx.$$

For n = 2, let $h(x_1, x_2)$ and $\eta(x_1, x_2)$ are real valued functions with support in $B_{\varepsilon}(0)$ such that

$$\int_{\mathbb{R}^2} h(x_1, x_2) \, dx_1 dx_2 = \int_{\mathbb{R}^2} \eta(x_1, x_2) \, dx_1 dx_2 = 1$$

For $\omega = h - \eta$,

$$\int_{\mathbb{R}^2} \omega(x) dx = 0$$

For $f(x_1, x_2) = \{h(A(x_1, x_2)) - \eta(A(x_1, x_2))\} \det J_A(x_1, x_2)$, we show there is an expression $g := g_1(x_1, x_2) dx_1 + g_2(x_1, x_2) dx_2$,

with support in $B_{\varepsilon}(0)$ such that

$$dg(x_1, x_2) = f(x_1, x_2) \, dx_1 dx_2$$

If we are able to show that, then

$$\int_{\mathbb{R}^2} f(x_1, x_2) \, dx_1 dx_2 = \int_{\mathbb{R}^2} dg(x_1, x_2) = 0.$$

as g, ω have compact support. Define g,

$$g(x_1, x_2) = \int_{-\infty}^{x_1} \left\{ f(t, x_2) - \left(\int_{-\infty}^{\infty} f(x_1, x_2) \, dx_1 \right) \tau(x_1) \right\} dt$$

where $\tau(x_1)$ is a function with the property

$$\int_{-\infty}^{\infty} \tau\left(x_1\right) dx_1 = 1$$

It is simply seen that $g(\infty, y) = g(x, \infty) = 0$. We have

$$\frac{\partial g_1}{\partial x_1} + \left(\int_{-\infty}^{\infty} f(x_1, x_2) \, dx_1 \right) \tau(x_1) = f(x_1, x_2)$$

Let us denote $g_3(x_2)$ by

$$g_3(x_2) = \int_{-\infty}^{\infty} f(x_1, x_2) \, dx_1$$

This is a function of a single variable x_2 and thus, there is $g_4(x_2)$ such that

$$g_3\left(x_2\right) = dg_4\left(x_2\right),$$

and therefore

$$f(x_1, x_2) = \frac{\partial g_1}{\partial x_1} + \frac{\partial \left(g_4(x_2) \tau(x_1)\right)}{\partial x_2}$$

Now assume for integral is independent of h for n = N. This means in \mathbb{R}^N there exists an $\omega = f(x)dx$ such that $\int_{\mathbb{R}^N} \omega(x)dx = 0$ with support in some $B_{\varepsilon}(0)$ and there exists a g with support in $B_{\varepsilon}(0)$ such that $\omega = d(g)$.

We show the property is true for n = N + 1. Let $x_1 = t$, $(t, x) = (t, x_2 \dots x_{N+1})$. Consider,

$$g(t,x) = \int_{-\infty}^{t} (f(s,x) - \tau(s)r(x))ds$$

where $\tau(s)$ has the property,

$$\int_{-\infty}^{\infty} \tau(t) dt = 1$$

and set

$$r(x) = \int_{-\infty}^{\infty} f(t, x) dt$$

This is a function of N dimensions. By induction hypothesis, $\int r(x)dx = 0$ and there exist g_1, \ldots, g_N such that

$$r(x) = \sum_{j=1}^{N} \frac{\partial g_j}{\partial x_j}$$

and each g_j are supported in $B_{\varepsilon}(0)$. Now

$$g(t,x) = \int_{-\infty}^{t} \left(f(s,x) - \tau(s) \sum_{j=1}^{N} \frac{\partial g_j}{\partial x_j} \right) ds$$

This integral vanishes in t as g has support in $B_{\varepsilon}(0)$, Thus

$$\frac{\partial g(t,x)}{\partial x_t} = f(t,x) - \tau(t)r(x)$$

which gives us

$$f(t,x) = \frac{\partial g(t,x)}{\partial x_t} + \sum_{j=1}^N \frac{\partial (g_j(x)\tau(t))}{\partial x_j}$$

Assume $D \subset \mathbb{R}^N$ is a bounded, open set and $f : D \to \mathbb{R}$ is C^1 . Now suppose $f : \overline{D} \to \mathbb{R}^N$ is smooth and $0 \notin f(\partial D)$. We call a point **regular** if $J_f(x) \neq 0$ whenever $x \in D$ and f(x) = 0.

Lemma 3. We claim that the set of regular points $f^{-1}(\{0\}) = \{x_1, x_2, \dots, x_N\}$ is finite.

Proof: Note that $f^{-1}(\{0\})$ is compact since it is a closed subset of \overline{D} . Now for $\{x_i\}_{i\in\mathbb{N}} \in f^{-1}(\{0\})$, the Inverse Function Theorem ensures that we can find $\epsilon > 0 \ni f^{-1}(B_{\epsilon}(0))$ is a union of disjoint neighbourhoods say $B_{\epsilon_i}(x_i)$ of $x_i \ni$ each $B_{\epsilon_i}(x_i) \cap D = x_i$. Therefore,

$$f^{-1}(\{0\}) \subset \bigcup_{x_i \in f^{-1}(\{0\})} B_{\epsilon_i}(x_i), \quad i \in \mathbb{N}, \epsilon_i > 0$$

Since $\{B_{\epsilon_i}(x_i)\}$ is an open cover of $f^{-1}(\{0\})$ we can find a finite subcover \ni ,

$$f^{-1}(\{0\}) \in B_{\epsilon_1}(x_1) \cup B_{\epsilon_2}(x_2) \cup \dots B_{\epsilon_i}(x_n) \implies f^{-1}(\{0\})$$
 is finite.

Proposition 1. Degree is Homotopically invariant: Let $f: \overline{D} \subset \mathbb{R}^N \to \mathbb{R}^N$ be an open bounded set and f_t be defined as $f_t(x): [0,1] \times \overline{D} \to \mathbb{R}^N$ such that it is $C^0([0,1] \times \overline{D})$ and $C^1(D)$ for each $t \in [0,1]$. Suppose $\forall t, 0 \notin f_t(\partial D)$. Then $deg(f_t, D, 0)$ is independent of t.

Proof: Choose $h(x) : U \subset \mathbb{R}^N \to \mathbb{R}^N \ni \int_{\mathbb{R}^N} h(x) = 1$ with support in a small neighbourhood U of $0 \ni U \cap f_t(\partial D) = \phi$. It's clear from the definition of an integral and $f_t(x)$ that.

$$deg(f_t, D, 0) = \int_D h(f_t(x)) \cdot |J_{ft}(x)| dx$$

is continous. Since $deg(f_t, D, 0)$ is an integrer, it doesn't depend on t, we get $deg(f, D, 0) = deg(f_t, D, 0)$. \Box

Remark 2. I now show that,

$$\sum_{x_i \in f^{-1}(0)} \operatorname{sign}(|J_f(x_i)|) = \int_D h(f(x)) |J_f(x)| dx$$

where h, f, D are defined as before.

We know that f(x) has a finite number of regular zero points in D. Therefore we can choose the $h \ni$ it's support is in $B_{\epsilon}(0) \subseteq \bigcap f(B_{\epsilon i}(x_i))$ (By the inverse function theorem). Since $J_f(x_i)$ has a fixed sign in each $B_{\epsilon_i}(x_i)$,

$$\begin{split} \int_{B_{\epsilon i}(x_i)} h(f(x)) |J_f(x)| dx &= \operatorname{sign}(|J_f(x_i)|) \int_{B_{\epsilon i}(x_i)} h(f(x)) dx, \\ &= \operatorname{sign}(|J_f(x_i)|) \int_{f(B_{\epsilon i}(x_i))} h(f(y)) dy \\ &= \operatorname{sign}(|J_f(x_i)|) \end{split}$$

Therefore, (here n is the number of regular points)

$$\int_D h(f(x))|J_f(x)|dx = \sum_{i=1}^n \int_{B_{\epsilon i}(x_i)} h(f(x))dx$$
$$= \sum_{i=1}^n \operatorname{sign} |J_f(x_i)|$$
$$= deg(f, D, 0)$$

(12)

We get (12) as the integral is 0 outside $\bigcup_{i=1}^{n} B_{\epsilon_i}(x_i)$.

Lemma 4. Let $f \in C^1(U) \cap C(\overline{D})$, D is the open unit ball and $f(x).x > 0 \forall x \in \partial D$. Show $\exists c \in D \ni f(c) = 0$.

Proof: Well we know $f(x) \neq 0$ on ∂D . Define the homotopy $f_t(x) : [0,1] \times \overline{D} \to \mathbb{R}^N$, $f_t(x) = tx + (1-x)f(x) \quad \forall x \in \overline{D}$. Since neither f(x) nor $x \neq 0$. This gives us $f_t(x) \neq 0$ for $x \in \partial D$ as $f_0(x) \neq 0$ and $f_1(x) \neq 0$ for $x \in \partial D \implies f_t(x).x > (1-t)x.x > 0$ for $x \in \partial D$ and $t \in (0,1)$. Therefore, $\deg(f, D, 0) = \deg(f_t, D, 0) = \deg(I, D, 0) = 1 \neq 0$. Thus $\exists c \ni f(c) = 0$.

Theorem 2. Brouwer Fixed Point Theorem: Let $f \in C^1(D) \cap C(\overline{D}), f : D \to D \implies \exists c \in D \ni f(c) = c.$

Proof: Let $D = B_{\epsilon}(0), \epsilon > 0$ and define the continuous homotopy $f_t : [0,1] \times \overline{D} \to \mathbb{R}^N, f_t(x) = x - tf(x)$. Assume $f(x) \neq x$ for $x \in \partial D$, otherwise we are done. Then if $x \in \partial D$ and $t \in [0,1)$, if x - tf(x) = 0then $\epsilon = ||x|| = t||f(x)|| < ||f(x)|| \le \epsilon$ a contradiction and if $x \in \partial D$ and t = 1, x - tf(x) = 0 implies $\epsilon = ||x|| = ||f(x)||$, but we assumed $f(x) \neq x, \forall x \in \partial D$ again a contradiction. From this we get that the homotopy is defined and $f_t(x) \neq 0$ for $x \in \partial D$. By previous lemma, $\deg(I - f(x), D, 0) = \deg(I, D, 0) = 1.\Box$

5. Infinite dimensional case

Consider the l^2 -space

$$l^{2} = \left\{ (x_{1}, x_{2}, \cdots) \mid ||x|| = \left| \sum_{i=1}^{\infty} x_{i}^{2} \right|^{\frac{1}{2}} < \infty \right\}.$$

This is a metric space with metric

$$d(x_i, y_i) = \sqrt{\sum_{i=1}^{\infty} (x_i - y_i)^2}.$$

Let B be the unit ball.

$$f: \bar{B}_{l^2} \to \bar{B}_{l^2}, \ (x_1, x_2, \cdots) \mapsto \left(\sqrt{1 - \|x\|^2}, x_1, x_2, \cdots\right).$$

The above mapping is continuous as for when $x, y \in l^2, d(x, y) \to 0$ we have,

$$\begin{split} \|f(x) - f(y)\| &= \sqrt{\left|\sqrt{1 - \|x\|^2} - \sqrt{1 - \|y\|^2}\right|^2} + \|x - y\|^2 \\ &\leq \sqrt{\|\|x\|^2 - \|y\|^2} + \|x - y\|^2 \\ &\leq \sqrt{(\|x\| + \|y\|)\|x - y\|} + \|x - y\|^2 \\ &\leq \sqrt{2\|x - y\|} + \|x - y\|^2} \to 0. \end{split}$$

However f doesn't have a fixed point. Assume f does, then since ||f(x)|| = 1 we get ||x|| = 1. Since $x = f(x) \implies x = \left(\sqrt{1 - ||x||^2}, x_1, x_2, \cdots\right) \implies x = (0, x_1, x_2, \ldots) \implies x_1 = 0, x_2 = x_1, x_3 = x_2 \ldots \implies x_1 = 0 = x_2 = \ldots$. But ||x|| = 1. A contradiction.

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